

# Simulation Studies on Re-emission and Thermal Desorption of Deuterium from a Tungsten Material

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## 1) Abstract

The hydrogen retention and recycling properties of first wall materials are important factors controlling the particle balance. We have tried to follow Garcia-Rosales's experiment [1] on fluxes on re-emission and desorption of deuterium from a tungsten material using the ACAT-DIFFUSE code[2].

[1] C. Garcia-Rosales et al., J. Nucl. Mater. **233-237** (1996) 803

[2] Y. Yamamura, Nucl. Instr. Meth. **B28** (1987) 17

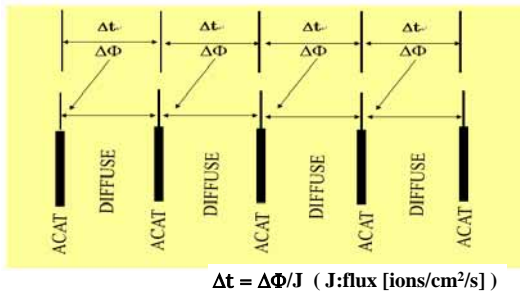
## 2) ACAT-DIFFUSE model

### Logical representation

$$\text{ACAT-DIFFUSE} = [(\text{ACAT}) + (\text{DIFFUSE})]^n$$

$$(n = \Phi / \Delta\Phi, \Phi: \text{Total fluence [ions/cm}^2\text{)})$$

### ACAT-DIFFUSE model



$$\Delta t = \Delta\Phi / J \quad (J: \text{flux [ions/cm}^2\text{/s)})$$

### 2-1) ACAT Routine

Monte Carlo method

Binary collision approximation

Amorphous target

Unit cell

### 2-2) DIFFUSE Routine

$$\frac{\partial c_j(x,t)}{\partial t} = \nabla [D_j(x,t) \nabla c_j(x,t)] + G_j(x,t_0) - \sum_{i=1}^2 \frac{\partial c_{Tj}^i(x,t)}{\partial t}$$

$j$ : j-atom of a multi-component material,  $c_j(x,t)$ : concentration of j-atom ( $\text{cm}^{-3}$ ),  $D_j$ : diffusion coefficient of j-atom ( $\text{cm}^2/\text{s}$ ),  $c_{Tj}^i$ : concentration of j-atom in i-defect ( $\text{cm}^{-3}$ ),  $G_j(x,t_0)$ : source term.

$$\frac{\partial c_{Tj}^i(x,t)}{\partial t} = \frac{D_j(x,t)c_j(x,t)C_{Te}^i(x,t)}{\lambda^2} - c_{Tj}^i(x,t)v_0 \exp(-E_T^i/kT)$$

$\lambda$ : the jump distance (cm),  $v_0$ : the detrapping attempt frequency ( $\text{s}^{-1}$ ),  $E_T^i$ : the detrapping energy of i-atom (eV).

$$C_{Te}^i(x,t) = C_{Te}^i(x,t_0) - \sum_j f_j^i c_j^i(x,t)$$

$f_j^i$ : the saturation of the i-defect due to j-atom.

$$\frac{\partial c_j}{\partial x} = \frac{K_r}{D_j} c_j^2(x=0), \quad [J = K_r c_j^2(x=0)]$$

$K_r$ : recombination coefficient ( $\text{cm}^3/\text{s}$ ),  $D_j$ : diffusion coefficient ( $\text{cm}^2/\text{s}$ ),  $c_j(x=0)$ : surface density of deuterium ( $\text{cm}^{-2}$ ).

$$K_r = \frac{K_0}{\sqrt{T}} \exp\left[-\frac{E_r}{kT}\right]$$

$K_0$ : preexponential factor ( $\text{cm} \cdot \text{K}^{1/2}/\text{s}$ ),  $E_r$ : activation energy for recombination (eV),  $T$ : Temperature (K).

## 4) Results

Fig.1 : TDS Simulation

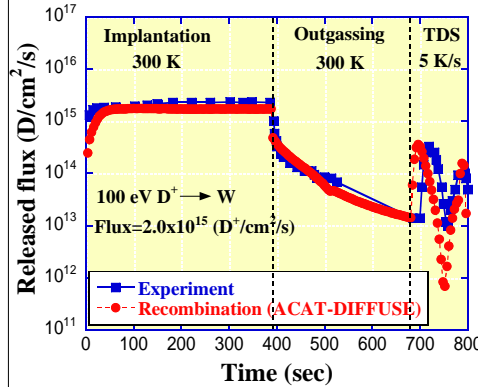


Fig.2 : Flux under the implantation

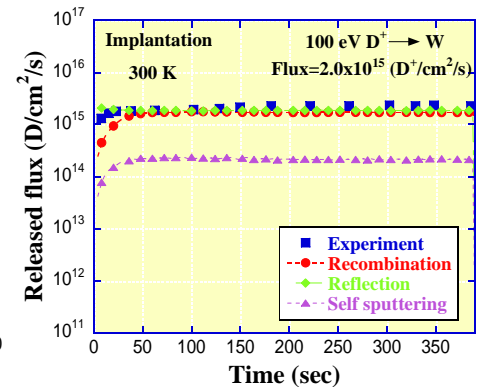


Fig.3 : Depth profile of D in C

(W.R.Wampler et al., JNM 103&104 (1981) 509)

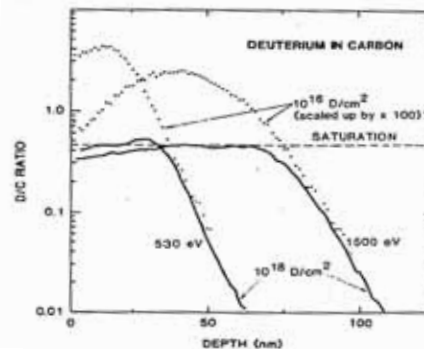


Fig.4 : Depth profile of D in W

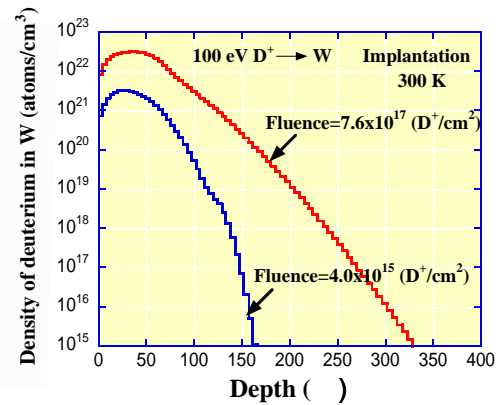


Fig. 5 : Mobile and trapped particles

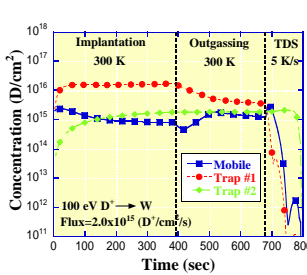
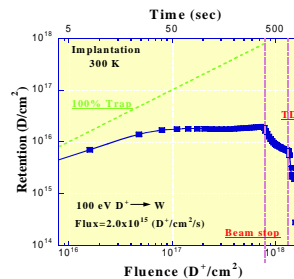


Fig.6 : Retention of D in W



### Parameters for the simulations

Parameters	Values
Diffusion	
$D_0$	$1.0 \times 10^{-8} \text{ cm}^2/\text{s}$
$E_D$	0.39 eV
Recombination	
$K_0$	$7.7 \times 10^{-3} \text{ cm} \cdot \text{K}^{1/2}/\text{s}$
$E_r$	-0.59 eV
Trap #1	
$E_{T,1}$	0.85 eV
$c_{T,1}$	0.05 traps/W
Trap #2	
$E_{T,2}$	2.2 eV
$c_{T,2}$	0.003 traps/W

## 5) Summary

➤ The ACAT-DIFFUSE code can predict the trend of the TDS experiment.

➤ In graphite almost all deuterium atoms are trapped in the implantation zone and don't diffuse, whereas in tungsten they distribute in a large zone larger than the implantation zone.