

Behavior of Reversible Hydrogen Getters under the Pulsed Plasma Heat Loads



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Motivation

Hydride-forming materials (getters) are considered as

Introduction

Hydrogen getters can be used in fusion devices for the creation of a controlled flow of hydrogen isotopes from the material surfaces to be exposed to hot plasma.

The paper presents results of investigations of hydrogen dynamics under the interaction of pulsed plasma streams with metal-hydrides on the base of Zr55V40Fe5 alloys in conditions of high-heat load.

RPI-IBIS coaxial multi-rod injector



Pulsed plasma streams were generated as a result of low-pressure high-current discharges performed between two coaxial sets of the multi-rod electrodes.
initial charging (supply) voltage was 30 kV

Initial charging (supply) voltage was 50 k v

- >delay time of the high-voltage pulse in relation to the injection of hydrogen into the inter-electrode volume was between 130-210 us
- > distance between the electrode ends and the investigated target was 100 mm
- ≻energy density in plasma stream 14-15 J/cm²

>energy of ions achieved 10 kV

≻pulse duration ~ 1-3 µs

rather prospective ones for the improvement of gas discharge cathodes. As the working gases in fusion devices hydrogen isotopes are usually used. Their additional flow, which should be matched with a pumping system, is ideal from the point of view of the providing the low-Z impurities, forming the shielding laver. So, this additional flow from constructional elements does not lead to any complication of the fusion reactor operation. It promotes the creation of a shielding gas target, which considerably screens the plasma energy density delivered to the material surface. One of the advantages of such a protection is that the shielding target is formed not from the target material (already eroded!) and not all time, but by the self-consistent way, i.e. only in the instants when plasma contacts with the material surface. The main aim of this work was to investigate shielding properties of reversible hydrogen getters under their irradiation with high-power plasma streams.

Power flux of the plasma stream used in our experiments was about 14 MW/cm², and similar one can be expected for off-normal events in ITER-like tokamaks. The pulse duration of order 1 µs is of course essentially shorter than at tokamak conditions. However, the first instants of the interaction are most important for an energy transfer to the material surface, because (due to the shielding) rather small part of plasma energy will be delivered later on.





Average values of N_e, T_e for different operational modes









Conclusions

>possibility of the effective shielding of material surfaces by desorbed hydrogen;

> the desorption above 2 x10¹⁹ particles/cm² achieved:

>the intermetallic phase of ZrCu₅ is formed during the crystallization;

≻the shielding layer formation with density higher than 10¹⁷ cm⁻³ (close to the sample surface) during parts of microsecond;

≻a difference in the delivered energy density, as registered for copper and metal hydride targets within the plateau region, can be explained by an influence of the destruction of the high-temperature hydrides ZrH₂ and ZrV₂H_{0.60};

>hydrogen is desorbed not only from a melted layer but also from a layer thicker than 20 µm, which is estimated to be heated up to such temperatures.

Discharge traces and temporal evolution of the H_{α} spectral line from the shielding layer, as recorded: a – for a copper target; b, c and d – for the getter target during the first, fourth and fifth plasma pulse, accordingly.





С

> the temporal evolution of the H_ spectral line from the near-surface plasma layer during the irradiation of copper and getter targets.

> the first peak in the H_{α} curve corresponds to the radiation of a free plasma stream and it can be observed even without the target. The second peak can be attributed to the radiation of the near-surface plasma layer (fig.A).

≻interaction of a plasma stream with the saturated getter target is accompanied by a significant increase in the intensity of H_ in the shielding layer formed due to the desorption of previously accumulated hydrogen(fig.B)

 $\succ intensity$ of H_{α} decreases with following pulses, but the duration of a H₋ signal is increased considerably, and it achieves 14-15 µs (fig.C-D).

Α







